Quantum Systems, Information, and Entanglement Lecture 1. Introduction to Quantum Systems

Spyros Tserkis

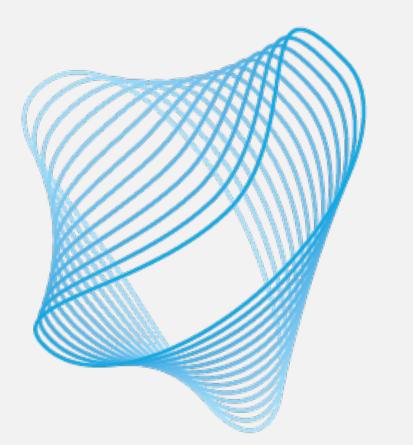
Postdoctoral Fellow

April 13, 2022









Center for Quantum Networks

There are four main thrusts to CQN:

Thrust 1: Quantum network architecture

Thrust 2: Quantum sub-system technologies

Thrust 3: Quantum materials, devices and fundamentals

Thrust 4: Societal impact of the Quantum Internet

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Center Director Professor, University of Arizona



Co-Deputy Director & Trust 4 Co-Professor of Law, University of



CQN Administrative Director



Innovation Ecosystem Director Executive in Residence, University of



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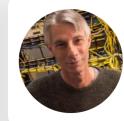


Testbed Co-Lead Assistant Professor, University o



Nena Bloom

CQN Evaluation Coordinator Evaluation Coordinator, Center for Science Teaching and Learning Northern Arizona University



Don Towsley

Thrust 1 Co-Lead Distinguished University Professor. **UMass Amherst**



Leandros Tassiulas

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Hong Tang

Thrust 2 Co-Lead Llewellyn West Jones, Jr. Professor of Electrical Engineering, Applied Physics, Yale



Mikhail Lukin

Thrust 2 Co-Lead George Vasmer Leverett Professor of Physics, Harvard



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Thrust 3 Co-Lead Assistant Professor, Harvard



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Thrust 3 Co-Lead Tiantsai Lin Professor of Electrical Engineering and Applied Physics,



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Thrust 4 Co-Lead iSchool Director and Associate Professor, University of Arizona



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Brian Smith

CQN Senior Personnel Associate Professor, University of





Course Outline

● Lecture 1 — Introduction to Quantum Systems (April 13, 2022)

● Lecture 2 — Teleportation and Entanglement (April 20, 2022)

● Lecture 3 — Decoherence and Quantum Networks (April 27, 2022)

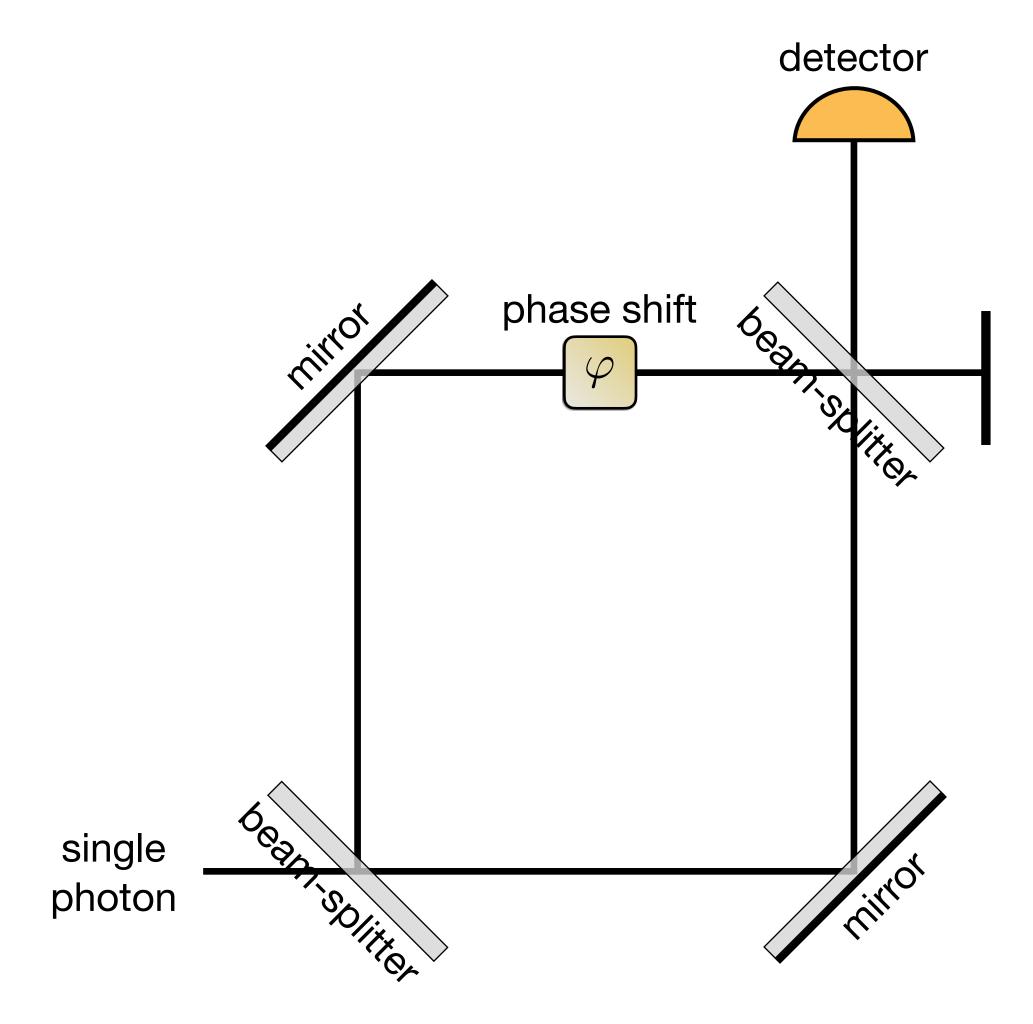
Lecture 1 — Introduction to Quantum Systems

- Introduction
 - When Classical Mechanics Fails
 - * Review of Classical Systems
- Quantum Observables
- Quantum States
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When Classical Mechanics Fails



Mach-Zehnder Interferometer

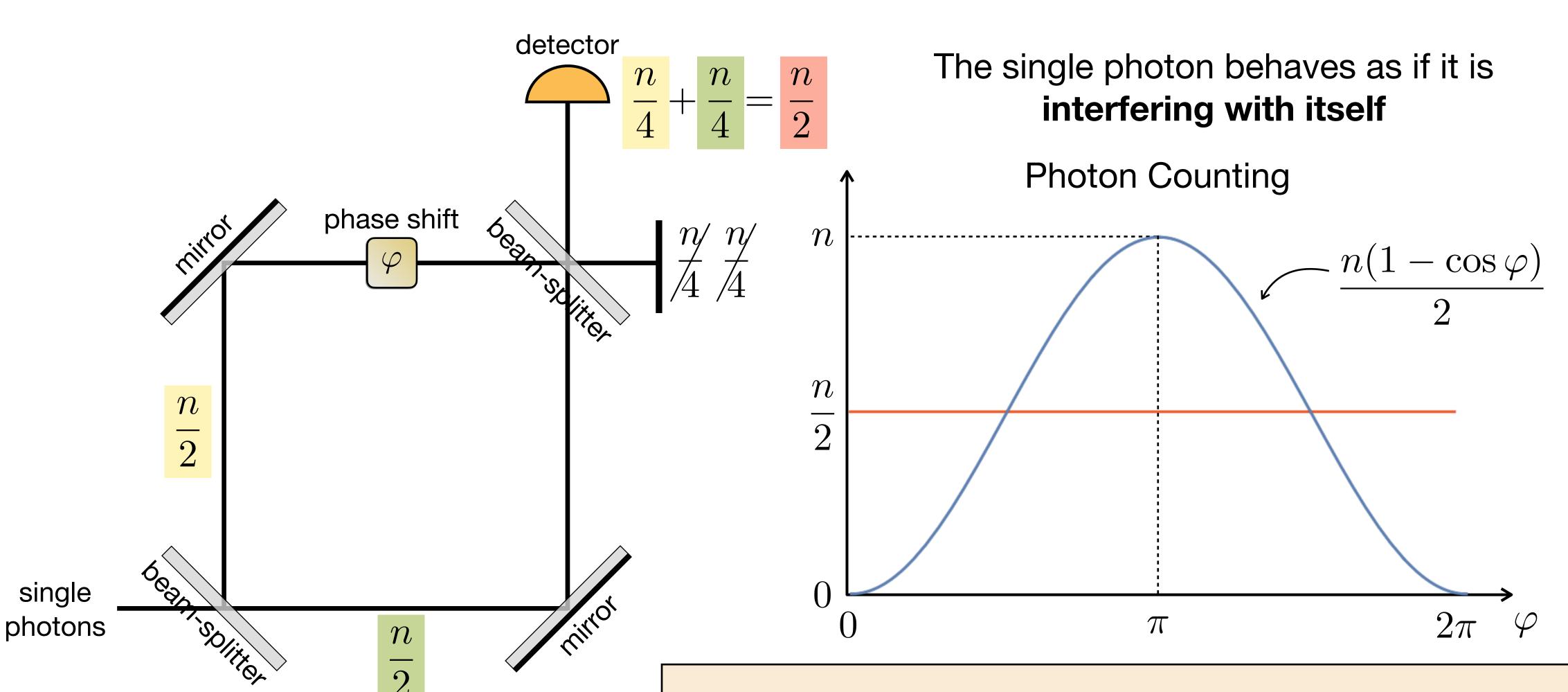
$$\varphi = QQ$$

$$\mathbb{P}(\text{red}) = 1/2$$

$$\mathbb{P}(\text{blue}) = 1/2$$

A photon takes each exit path with equal probability

When Classical Mechanics Fails



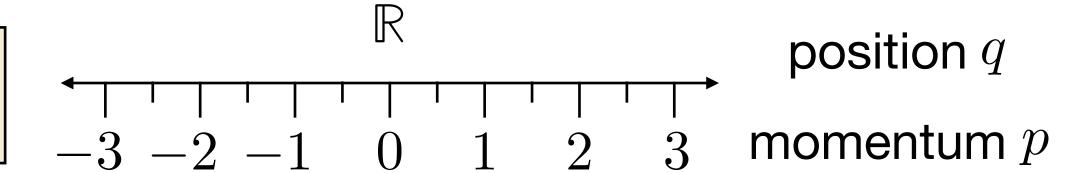
Mach-Zehnder Interferometer

Wave-particle Duality

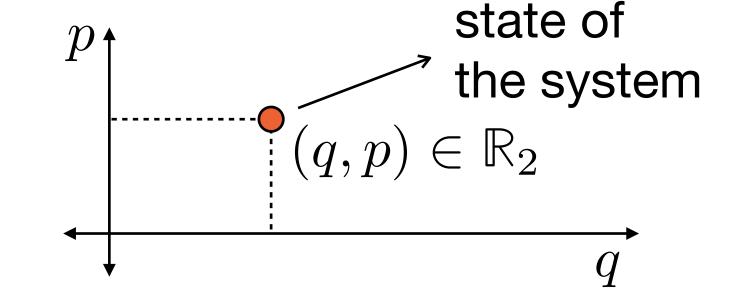
A particle can be thought of as both a particle and a wave

Classical Systems

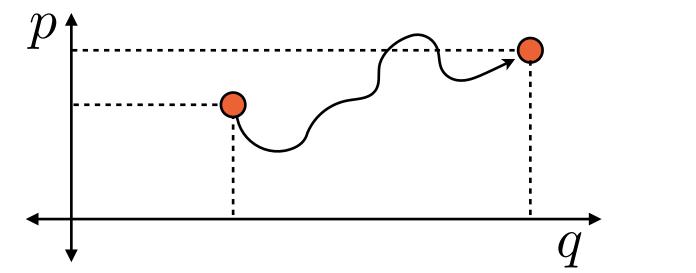
A property is represented by a real-valued number



A system is represented by a point in the phase space



Evolution of a property/state is deterministic



$$\frac{\mathrm{d}q}{\mathrm{d}t} = \frac{p}{m}$$

$$\frac{\mathrm{d}p}{\mathrm{d}t} = f$$

Dirac Notation

- A column vector is represented with a "ket", e.g., $|x\rangle = \begin{bmatrix} a \\ b \end{bmatrix}$
- A row vector is represented with a "bra", e.g., $\langle y|=\begin{bmatrix} c & d \end{bmatrix}$
- A ket can be transformed into a bra as follows: $|x\rangle \to \langle x| = \begin{bmatrix} a^* & b^* \end{bmatrix} = |x\rangle^\dagger$ (conjugate transpose)
- The inner product is represented as a bra-ket $\langle y|x\rangle=\begin{bmatrix}c&d\end{bmatrix}\begin{bmatrix}a\\b\end{bmatrix}=a\cdot c+b\cdot d$
- The outer product as $|x\rangle\!\langle y| = \begin{bmatrix} a \\ b \end{bmatrix} \begin{bmatrix} c & d \end{bmatrix} = \begin{bmatrix} ac & ad \\ bc & bd \end{bmatrix}$ E.g., $|x\rangle\!\langle y| \cdot |z\rangle = |x\rangle\,\langle y|z\rangle = \langle y|z\rangle\,|x\rangle$
- We consider the **computational basis** $|0\rangle=\begin{bmatrix}1\\0\end{bmatrix}$ and $|1\rangle=\begin{bmatrix}0\\1\end{bmatrix}$

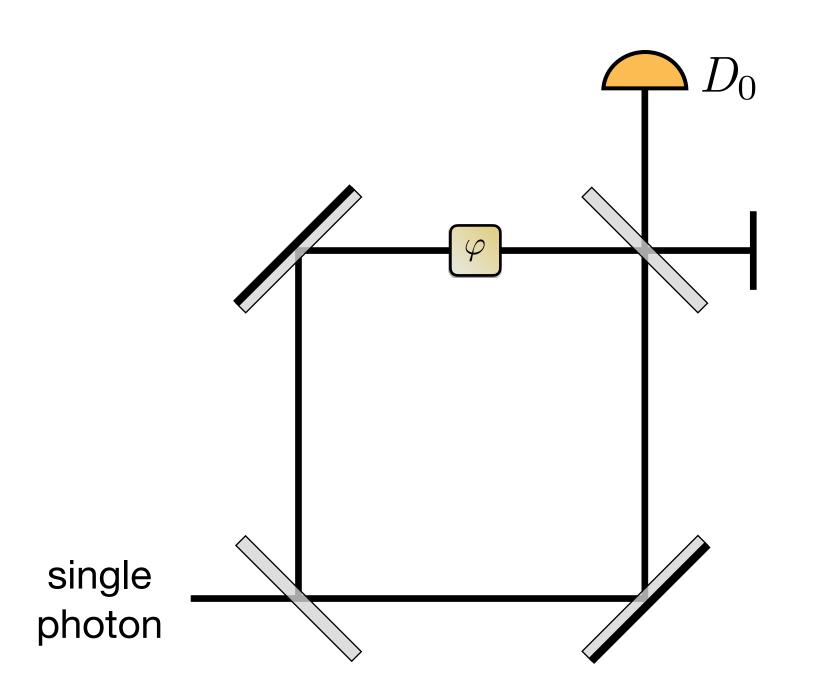
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A classical property is represented by a number

A quantum property, known as quantum observable, is represented by a matrix

- Those matrices have real eigenvalues and represent the possible outcomes of measurements.
- E.g., Photon detection is a binary property that can be represented by a matrix:



$$D_0=|0\rangle\!\langle 0|=\begin{bmatrix}1\\0\end{bmatrix}\begin{bmatrix}1&0\end{bmatrix}=\begin{bmatrix}1&0\\0&0\end{bmatrix}$$
 Eigenvalues: 0 and 1

Detection: 1

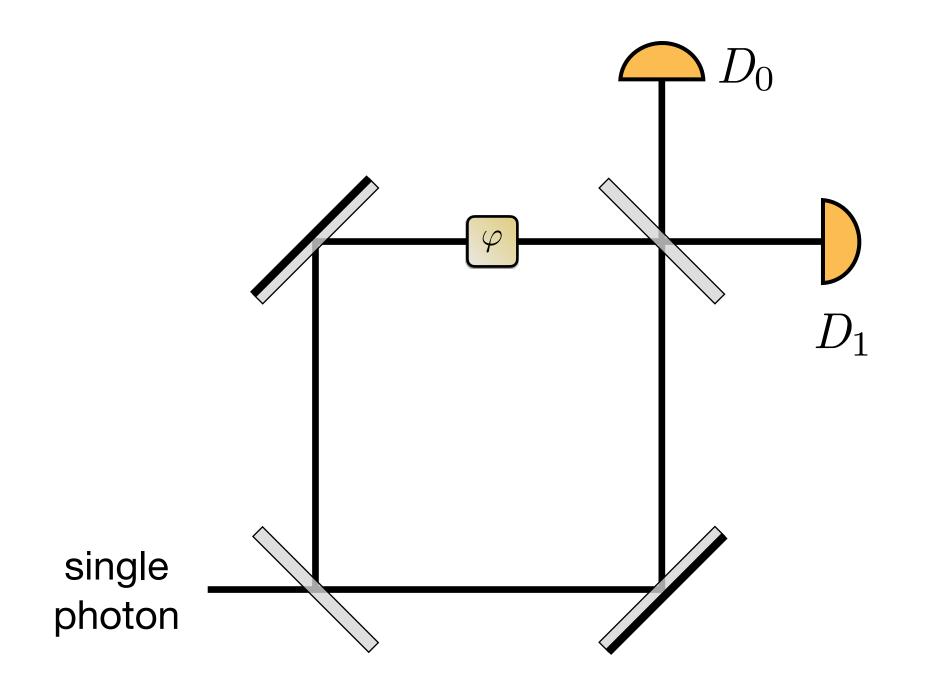
Non-detection: 0

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eigenvalues



$$D_0=|0\rangle\!\langle 0|=\begin{bmatrix}1\\0\end{bmatrix}\begin{bmatrix}1&0\end{bmatrix}=\begin{bmatrix}1&0\\0&0\end{bmatrix}$$
 Eigenvalues: 0 and 1

Detection: 1

Non-detection: 0

$$D_1$$
 $D_1=|1
angle\langle 1|=egin{bmatrix}0\\1\end{bmatrix}egin{bmatrix}0&1\end{bmatrix}=egin{bmatrix}0&0\\0&1\end{bmatrix}$ Eigenvalues: 0 and 1

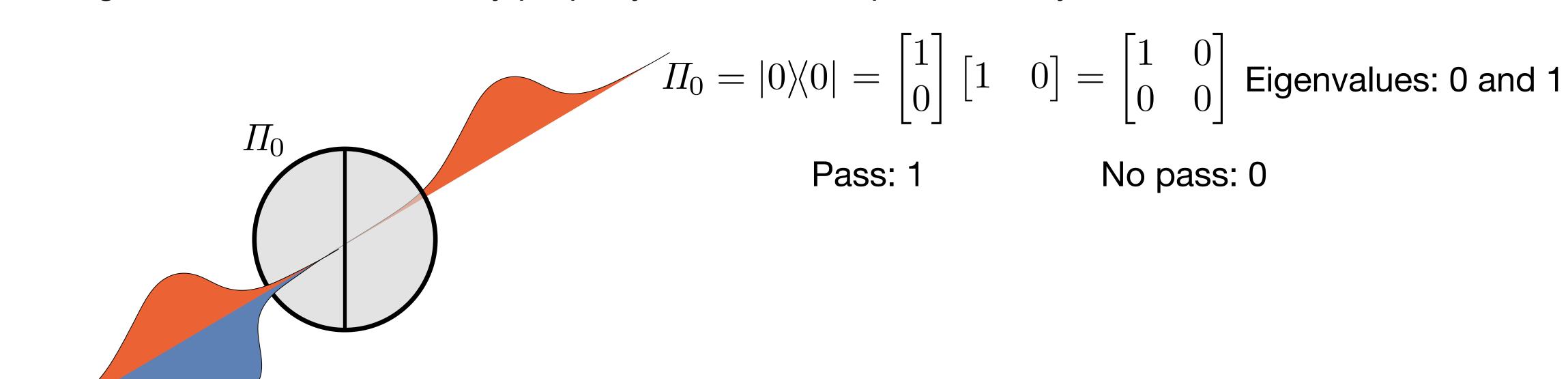
$$D = \lambda_0 D_0 + \lambda_1 D_1 = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \qquad \lambda_0 \neq \lambda_1 \qquad \lambda_1 = -1$$

quantum measurements

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A quantum property, known as quantum observable, is represented by a matrix

- Those matrices have real eigenvalues and represent the possible outcomes of measurements.
- E.g., Polarization is a binary property that can be represented by a matrix:

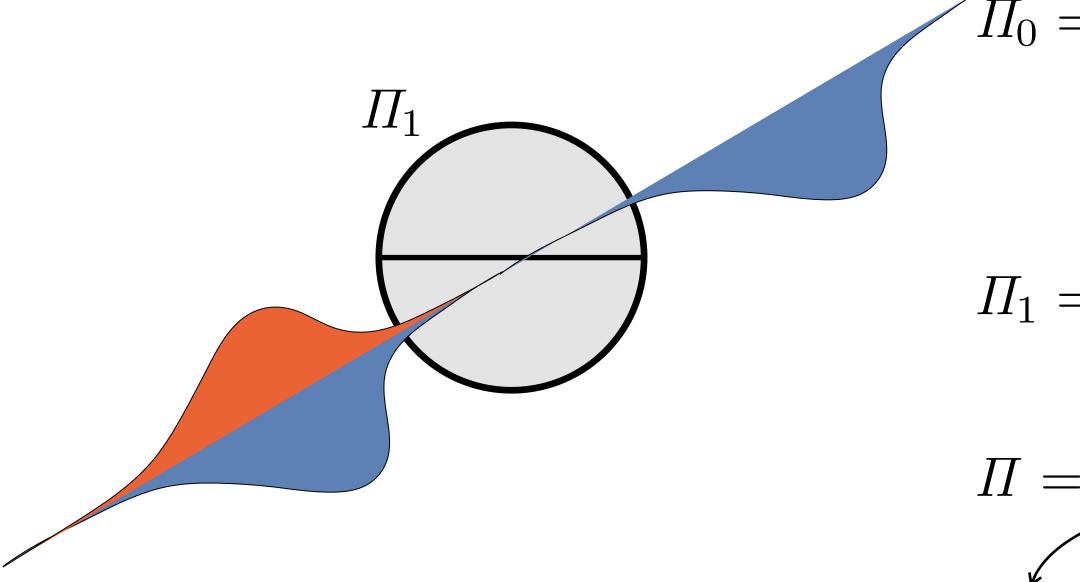


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A quantum property, known as quantum observable, is represented by a matrix

eigenvalues

- Those matrices have real eigenvalues and represent the possible outcomes of measurements.
- E.g., Polarization is a binary property that can be represented by a matrix:



$$I_0=|0
angle\langle 0|=egin{bmatrix}1\\0\end{bmatrix}\begin{bmatrix}1&0\end{bmatrix}=egin{bmatrix}1&0\\0&0\end{bmatrix}$$
 Eigenvalues: 0 and 1

Pass: 1

No pass: 0

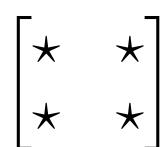
$$\varPi_1=|1\rangle\!\langle 1|=egin{bmatrix}0\\1\end{bmatrix}egin{bmatrix}0&1\end{bmatrix}=egin{bmatrix}0&0\\0&1\end{bmatrix}$$
 Eigenvalues: 0 and 1

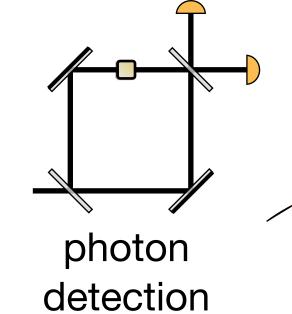
$$\Pi = \lambda_0 \Pi_0 + \lambda_1 \Pi_1 = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \qquad \lambda_0 \neq \lambda_1 \qquad \lambda_0 = 1$$
eigenvalues quantum measurements

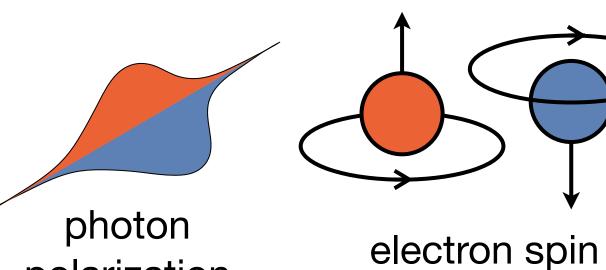
14

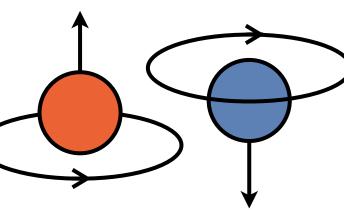
Different Types of Observables

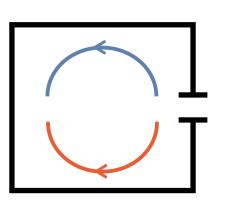
Observable with **two** possible values





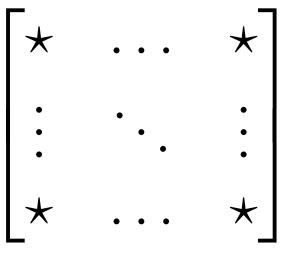






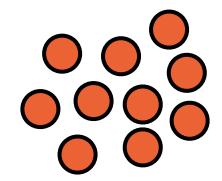
current flow

Observable with **finite** possible values



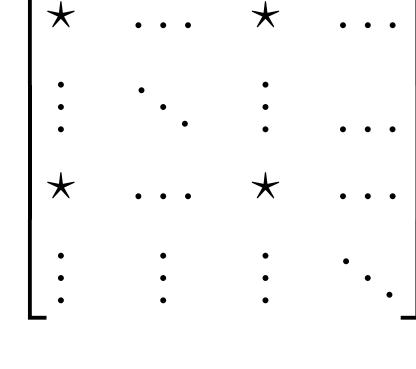
electron excitation

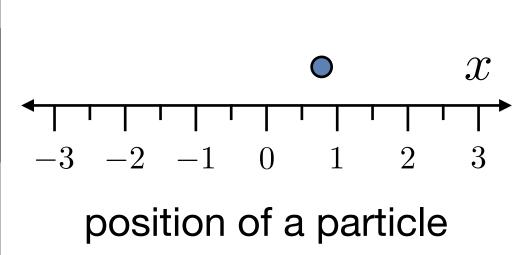
polarization

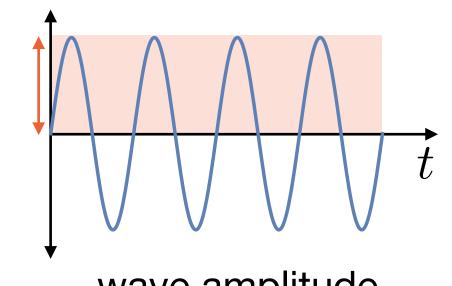


number of particles, e.g., photons

Observable with infinite possible values







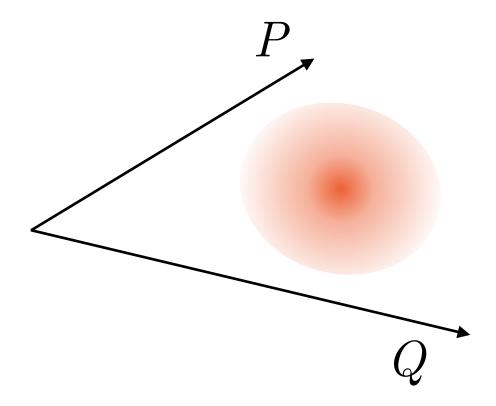
wave amplitude

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Quantum System

A classical system is represented by a point in the phase space

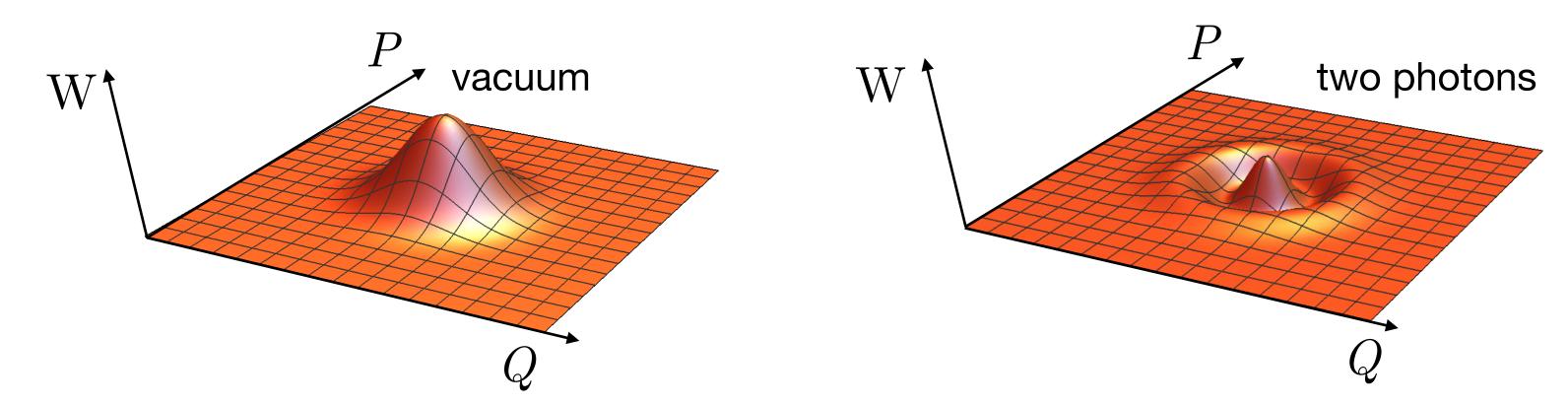


A quantum system cannot be represented as a point due to the uncertainty principle $\mathbb{V}(Q)\mathbb{V}(P)\geqslant \frac{\hbar^2}{4}$

Quantum System

A classical system is represented by a point in the phase space

A quantum system is represented by a function in the phase space, e.g., Wigner function



A quantum system cannot be represented as a point due to the **uncertainty principle** $\mathbb{V}(Q)\mathbb{V}(P)\geqslant \frac{\hbar^2}{4}$

$$W = \frac{1}{4\pi} \int_{-\infty}^{\infty} dx \exp\left\{\frac{ixp}{2}\right\} \left\langle q - \frac{x}{2} \middle| \psi \right\rangle \left\langle \psi \middle| q + \frac{x}{2} \right\rangle$$

q: outcome of Q

p: outcome of P

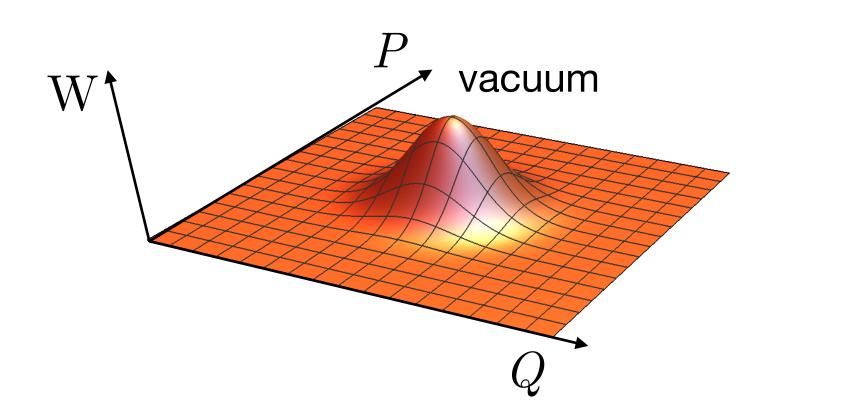
x: auxiliary variable

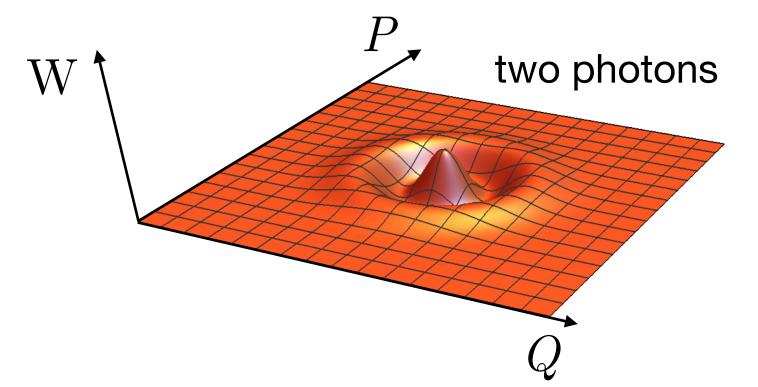
 $|\psi\rangle$: quantum state

Quantum System

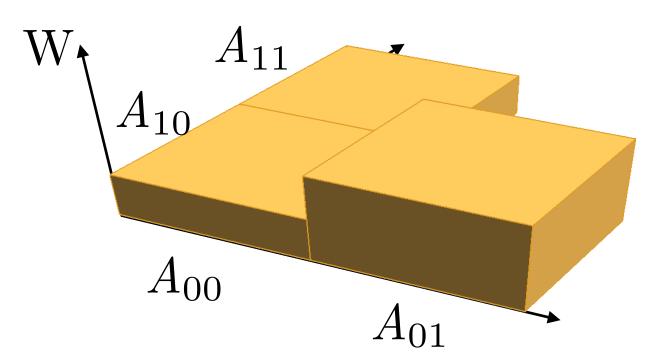
A classical system is represented by a point in the phase space

A quantum system is represented by a function in the phase space, e.g., Wigner function





A quantum system cannot be represented as a point due to the **uncertainty principle** $\mathbb{V}(Q)\mathbb{V}(P)\geqslant \frac{\hbar^2}{4}$



$$A_{ij} = \frac{1}{2} \left[(-1)^i Z + (-1)^j X + (-1)^{i+j} Y \right]$$

$$X = |0\rangle\langle 1| + |1\rangle\langle 0|$$

$$Y = -i(|0\rangle\langle 1| - |1\rangle\langle 0|)$$

$$Y = -i(|0\rangle\langle 1| - |1\rangle\langle 0|)$$

$$Z = |0\rangle\langle 0| - |1\rangle\langle 1|$$

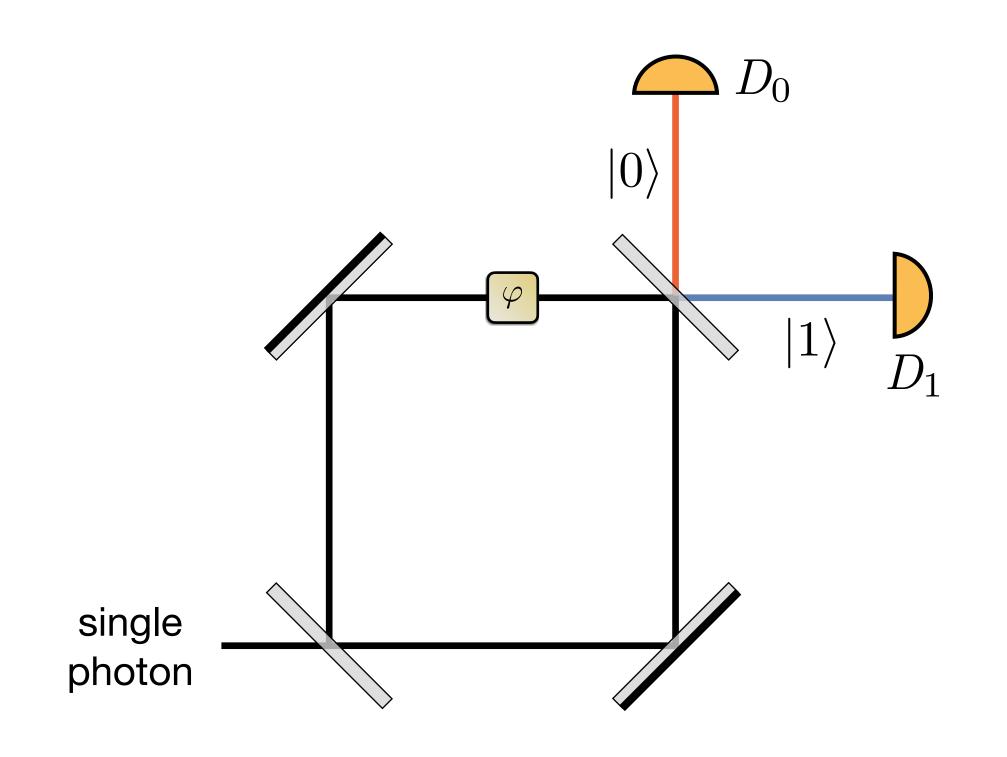
$$W_{ij} = \frac{1}{2} \operatorname{tr}(|\psi\rangle\langle\psi|A_{ij})$$

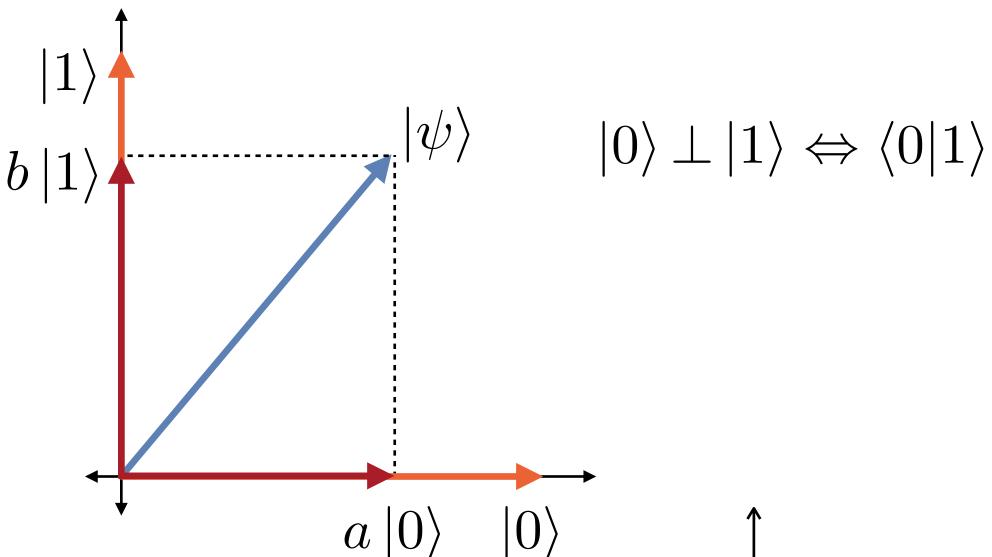
quantum state

Quantum State

A quantum system, known as quantum state, is represented by a vector

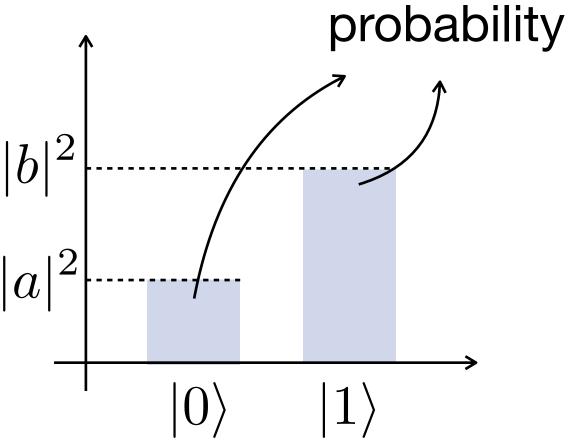
State Superposition: $|\psi\rangle=a\,|0\rangle+b\,|1\rangle$





All states have the same length, i.e., **normalized**

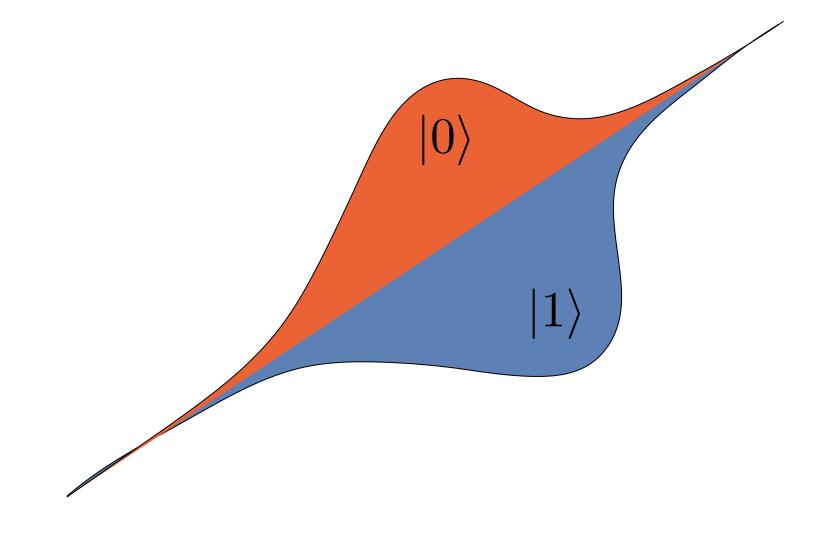
$$|a|^2 + |b|^2 = 1$$

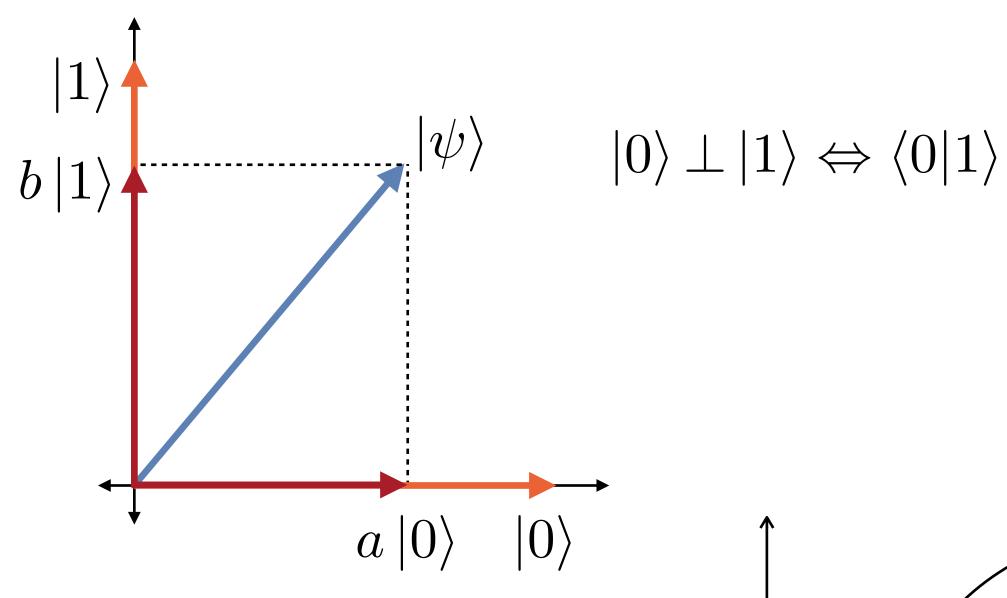


Quantum State

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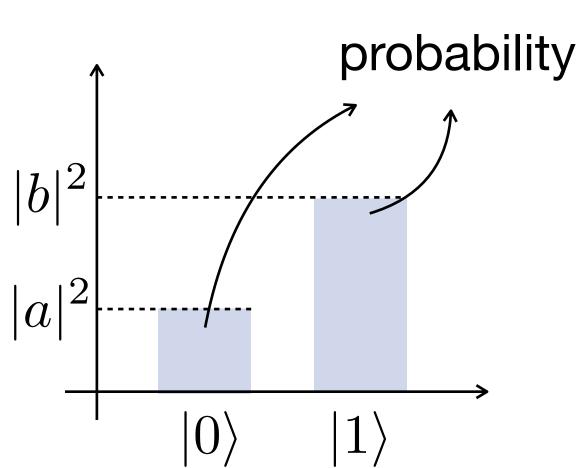
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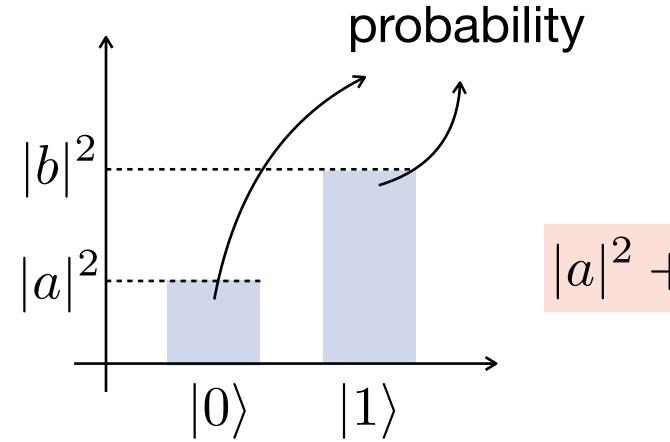
$$|a|^2 + |b|^2 = 1$$



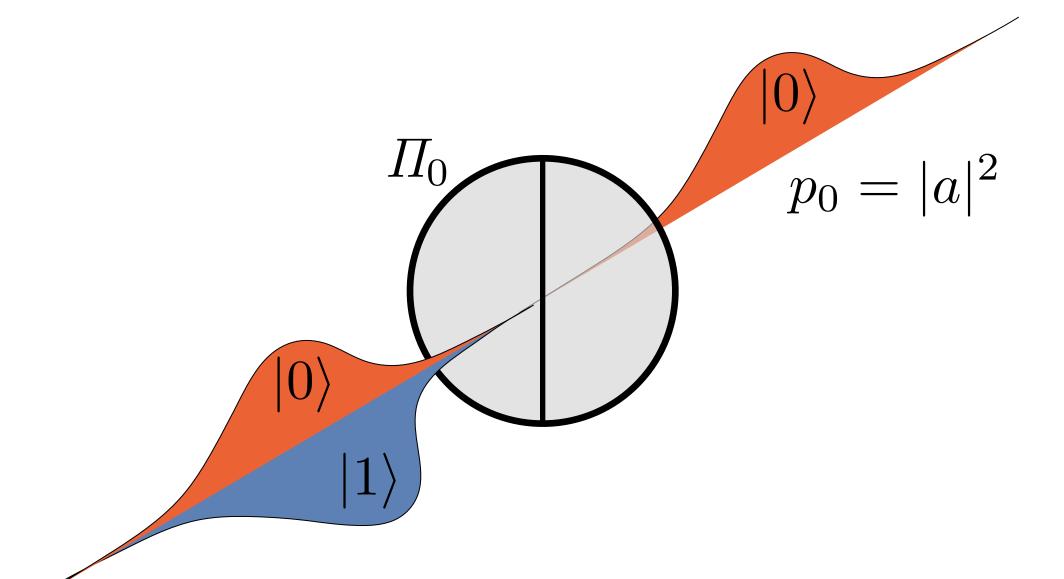
Quantum State after a Measurement

Let us have the quantum state $|\psi\rangle=a\,|0\rangle+b\,|1\rangle$

Post-measurement state $|\psi\rangle \to \frac{\varPi_i |\psi\rangle}{\sqrt{p_i}}$

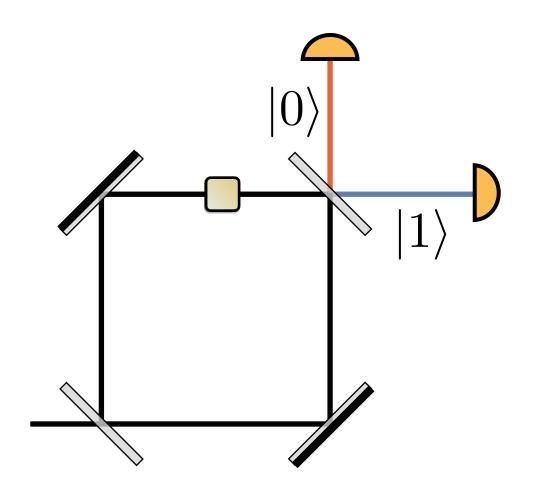


$$|a|^2 + |b|^2 = 1$$

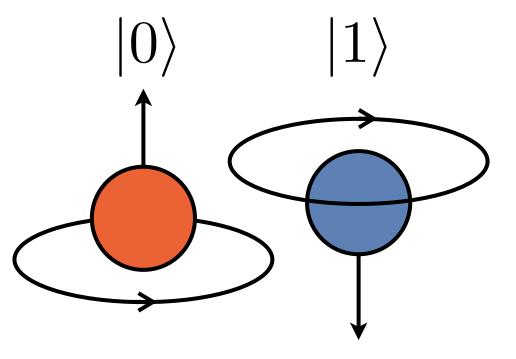


(collapse of the state) $p_1 = |b|^2$ $|1\rangle$

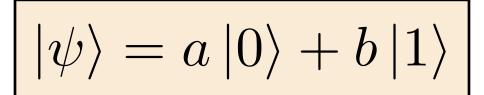
Examples of 2-dimensional States

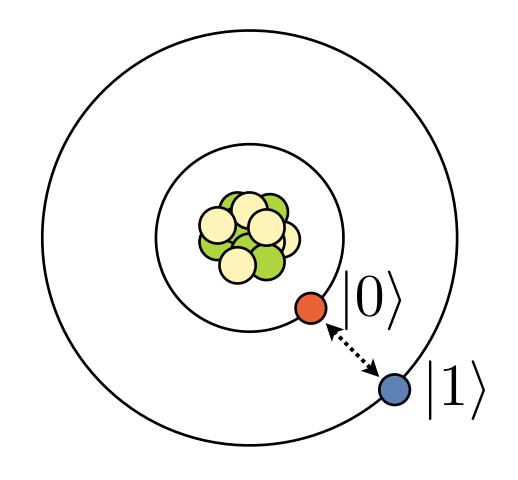


Photon Detection

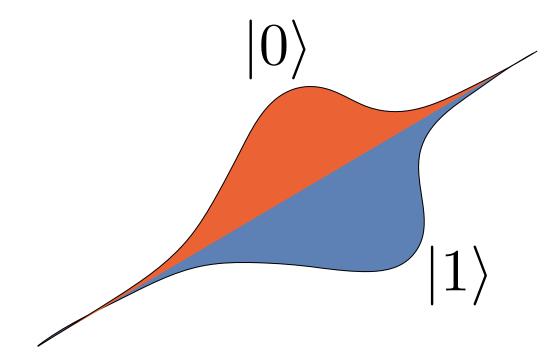


Electron Spin

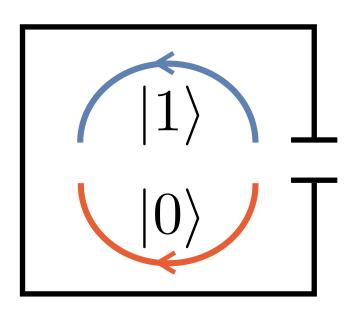




Electron Excitation



Photon Polarization

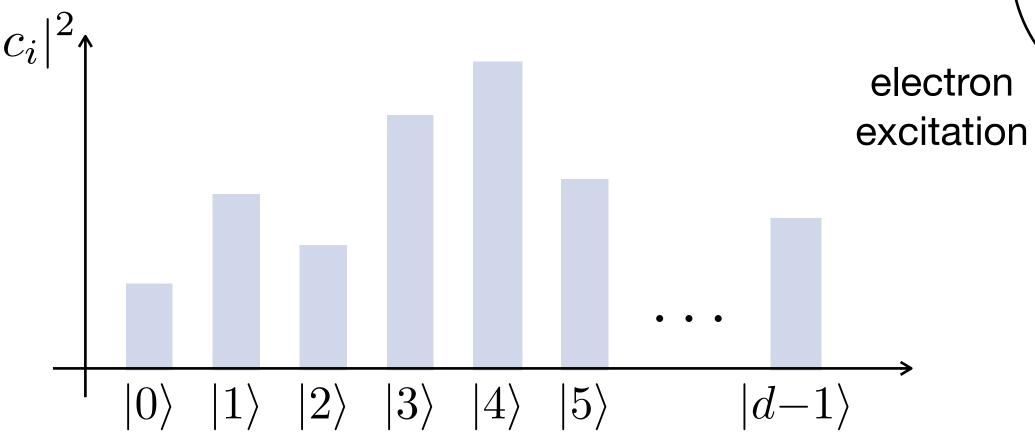


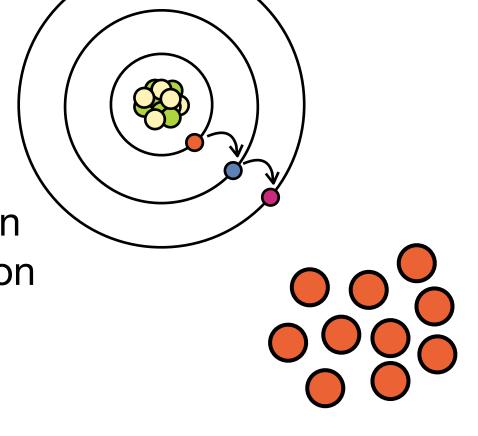
Current Flow

Higher-Dimensional States

A d-dimensional quantum state

$$|\psi\rangle = \begin{bmatrix} c_0 \\ c_1 \\ \vdots \\ c_{d-1} \end{bmatrix} = \sum_{i=0}^{d-1} c_i |i\rangle$$

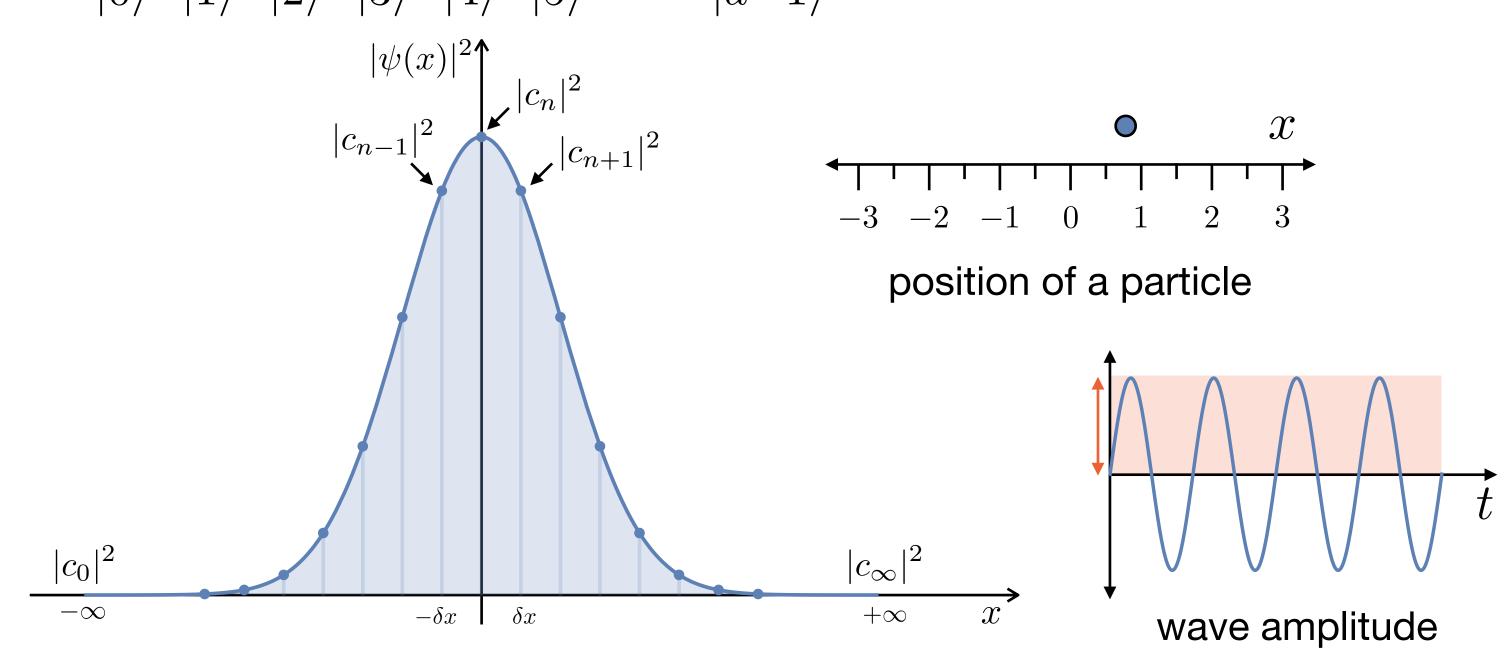




number of particles, e.g., photons

An infinite-dimensional state

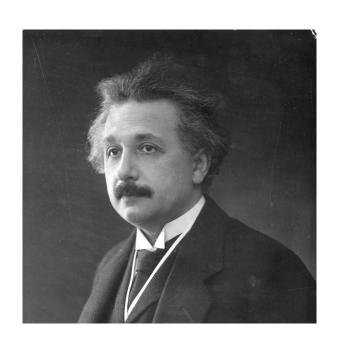
$$|\psi\rangle = \begin{bmatrix} c_0 \\ c_1 \\ \vdots \end{bmatrix} = \sum_{i=0}^{\infty} c_i |i\rangle$$



Quantum State Interpretation



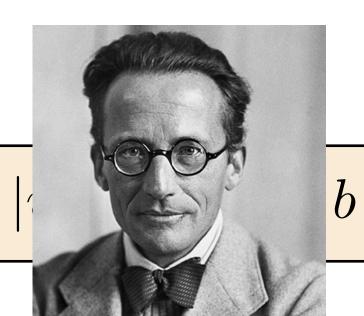




A. Einstein



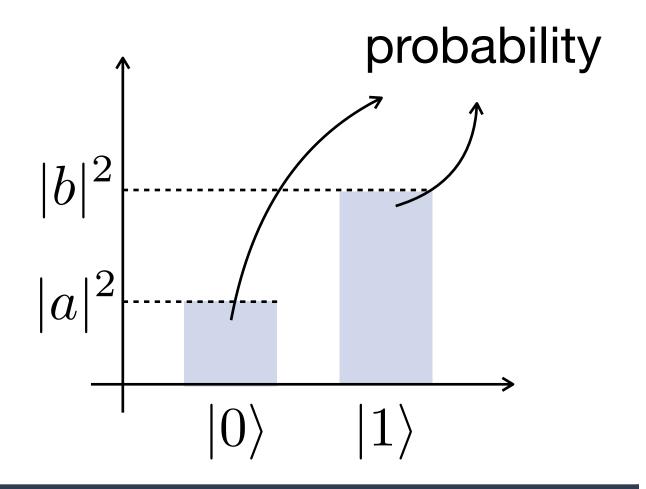
N. Bohr



E. Schrödinger



W. Heisenberg



Quantum State Interpretation

A quantum state corresponds to a **statistical ensemble** of independent and identically prepared copies of a quantum system

Probability Interpretation

Frequentism: The relative frequency of an event in the limit of sufficient many trials

A quantum state provides a complete description of an **individual** quantum system

Bayesianism: The degree of confidence of a hypothesis based on the prior knowledge

many worlds pilot waves

Copenhagen interpretation

quantum Bayesianism

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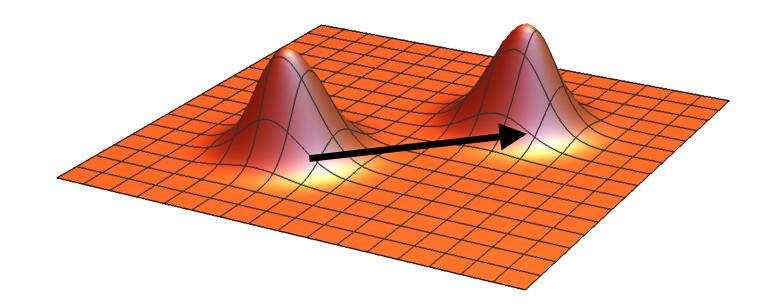
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Evolution of a Quantum System

The evolution of a classical property is deterministic

$$\frac{\mathrm{d}x}{\mathrm{d}t} = y$$

The evolution of a quantum property is deterministic $\frac{\mathrm{d}X}{\mathrm{d}t}=Y$



The evolution of a quantum state is described by a unitary transformation on the quantum state

$$|\psi\rangle \to U |\psi\rangle$$

- Unitary is a matrix that satisfies: $UU^\dagger = U^\dagger U = 1$ (preserves normalization)
- When $U = e^{iHt/\hbar}$ where H is the Hamiltonian of the quantum system and \hbar the reduced Planck constant the evolution is given by the **Schrödinger equation**

$$i\hbar \frac{\mathrm{d}}{\mathrm{d}t} |\psi\rangle = H |\psi\rangle$$

Example of Quantum Evolution

Let us have the quantum state $|\psi\rangle=\sqrt{\frac{1}{3}}\,|0\rangle+\sqrt{\frac{2}{3}}\,|1\rangle$

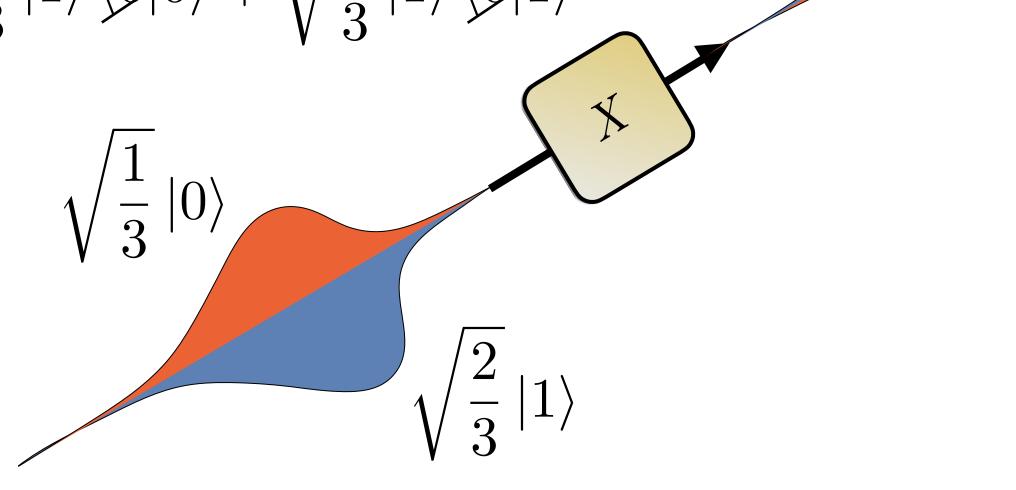
evolving according to the unitary matrix $X=|0\rangle\!\langle 1|+|1\rangle\!\langle 0|$

$$\begin{split} X \mid \psi \rangle &= (\mid 0 \rangle \langle 1 \mid + \mid 1 \rangle \langle 0 \mid) \left(\sqrt{\frac{1}{3}} \mid 0 \rangle + \sqrt{\frac{2}{3}} \mid 1 \rangle \right) \\ &= \sqrt{\frac{1}{3}} \mid 0 \rangle \langle 1 \mid 0 \rangle^0 + \sqrt{\frac{2}{3}} \mid 0 \rangle \langle 1 \mid 1 \rangle^1 + \sqrt{\frac{1}{3}} \mid 1 \rangle \langle 0 \mid 0 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^0 \\ &= \sqrt{\frac{1}{3}} \mid 0 \rangle \langle 1 \mid 0 \rangle^0 + \sqrt{\frac{2}{3}} \mid 0 \rangle \langle 1 \mid 1 \rangle^1 + \sqrt{\frac{1}{3}} \mid 1 \rangle \langle 0 \mid 0 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^0 \\ &= \sqrt{\frac{1}{3}} \mid 0 \rangle \langle 1 \mid 0 \rangle^0 + \sqrt{\frac{2}{3}} \mid 0 \rangle \langle 1 \mid 1 \rangle^1 + \sqrt{\frac{1}{3}} \mid 1 \rangle \langle 0 \mid 0 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^0 \\ &= \sqrt{\frac{1}{3}} \mid 0 \rangle \langle 1 \mid 0 \rangle^1 + \sqrt{\frac{2}{3}} \mid 0 \rangle \langle 1 \mid 1 \rangle^1 + \sqrt{\frac{1}{3}} \mid 1 \rangle \langle 0 \mid 0 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^0 \\ &= \sqrt{\frac{1}{3}} \mid 0 \rangle \langle 1 \mid 0 \rangle^1 + \sqrt{\frac{2}{3}} \mid 0 \rangle \langle 1 \mid 1 \rangle^1 + \sqrt{\frac{1}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle \langle 0 \mid 1 \rangle^1 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac{2}{3}} \mid 1 \rangle^2 \langle 0 \mid 1 \rangle^2 + \sqrt{\frac$$

$$= \sqrt{\frac{2}{3}} |0\rangle + \sqrt{\frac{1}{3}} |1\rangle$$

$$Z = |0\rangle\langle 0| - |1\rangle\langle 1| \qquad Z|\psi\rangle = 0$$

$$H = \frac{1}{\sqrt{2}}(X + Z) \qquad H |\psi\rangle = ?$$



 $\sqrt{\frac{2}{3}}\ket{0}$

Lecture 1 — Introduction to Quantum Systems

- Introduction
 - When Classical Mechanics Fails
 - * Review of Classical Systems
- Quantum Observables
- Quantum States
- Evolution in Quantum Systems
- Quantum Information

Quantum Information

In classical information theory we encode information onto bits

bit: 0 or 1

In quantum information theory we encode information onto quantum bits (qubits)

qubit:
$$a |0\rangle + b |1\rangle \rightarrow |0\rangle$$
 or $|1\rangle$

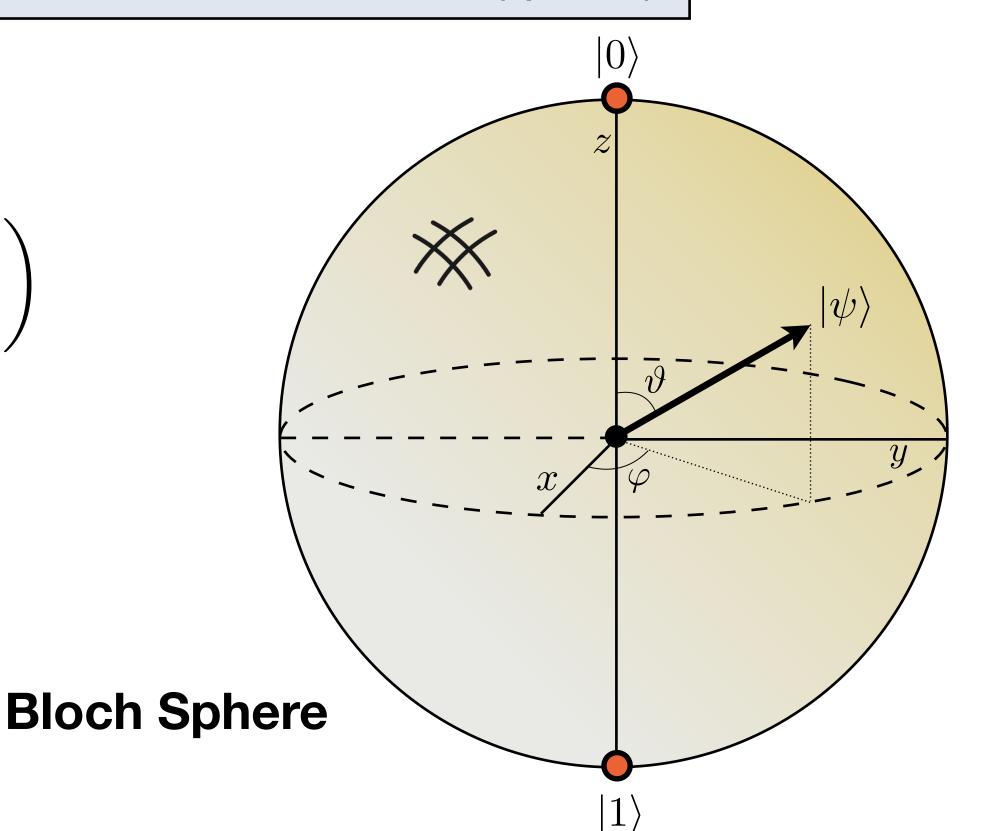
$$|\psi\rangle = a |0\rangle + b |1\rangle = e^{i\chi} \left(\cos\frac{\theta}{2} |0\rangle + e^{i\varphi} \sin\frac{\theta}{2} |1\rangle\right)$$

$$\chi \in [0, 2\pi]$$

$$\varphi \in [0, 2\pi]$$

$$\theta \in [0, \pi]$$

Qubit contains a bit as a special case



Next Week

● Lecture 1 — Introduction to Quantum Systems (April 13, 2022)

● Lecture 2 — Teleportation and Entanglement (April 20, 2022)

● Lecture 3 — Decoherence and Quantum Networks (April 27, 2022)

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Suggested Bibliography

- Quantum Mechanics
 - * J. Townsend "A Modern Approach to Quantum Mechanics"
 - * L. Ballentine "Quantum Mechanics"
- Quantum Information
 - * J. Audretsch "Entangled Systems"
 - * M. Nielsen, I. Chuang "Quantum Computation and Quantum Information"

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